

Ultrafast, broadband and tunable THz reflector based on high resistivity silicon.

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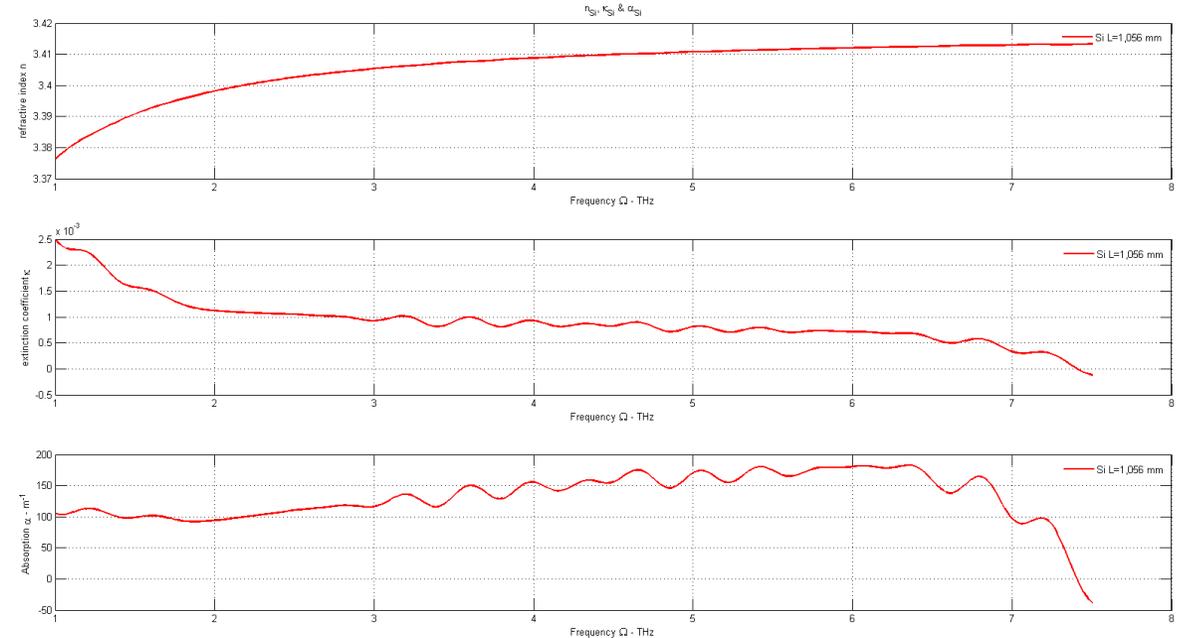
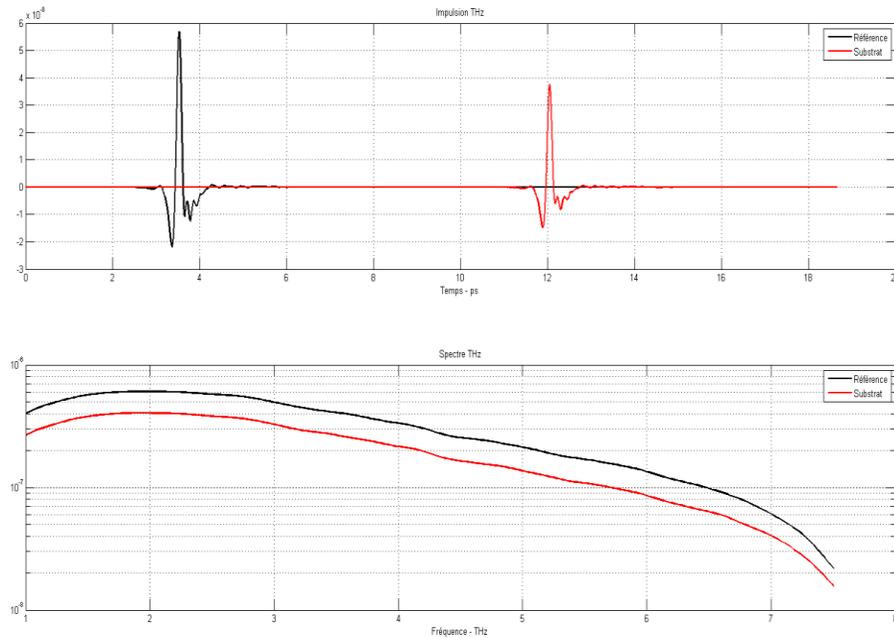
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Introduction

- Silicon is of utmost importance in electronic
- The control of defects and free carriers strongly impacts its performance.
- Nowadays, large silicon wafers with less than 10^{13} cm^{-3} doping concentration exhibit high resistivity $\sim 1000 \text{ } \Omega \cdot \text{cm}$.
- Such high resistivity silicon (HRSi) substrates have very low absorption: *below 3 THz a 1 mm thick HRSi substrate as an O.D. < 0.01.*
- They are often used as a low cost dichroic filter to reflect above Si bandgap radiation and transmit THz radiations.
- Hereafter we demonstrate they can also be used as ultrafast, broadband and tunable THz reflectors or density filters.

TDS-THz spectroscopy of the studied HRSi wafer

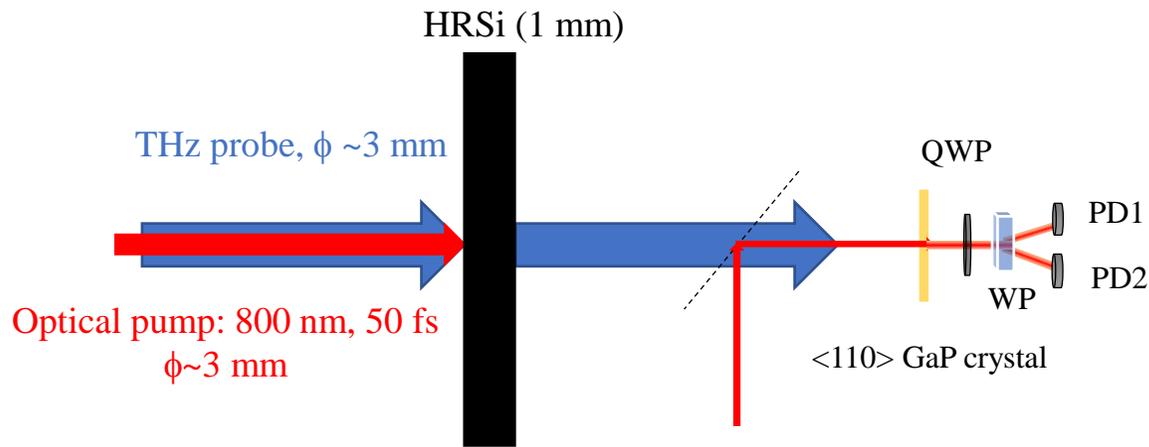
Experiment performed in transmission at normal incidence



- As the frequency increases, the index slightly increases from 3.378 (1 THz) to 3.412 (7 THz)
- Absorption is almost constant $\kappa \sim 0.01 \text{ cm}^{-1}$ from 1 THz to 7 THz
- The transmission at the THz peak is $T \approx 68\%$ in good agreement with the tabulated value $T \approx 70\%$

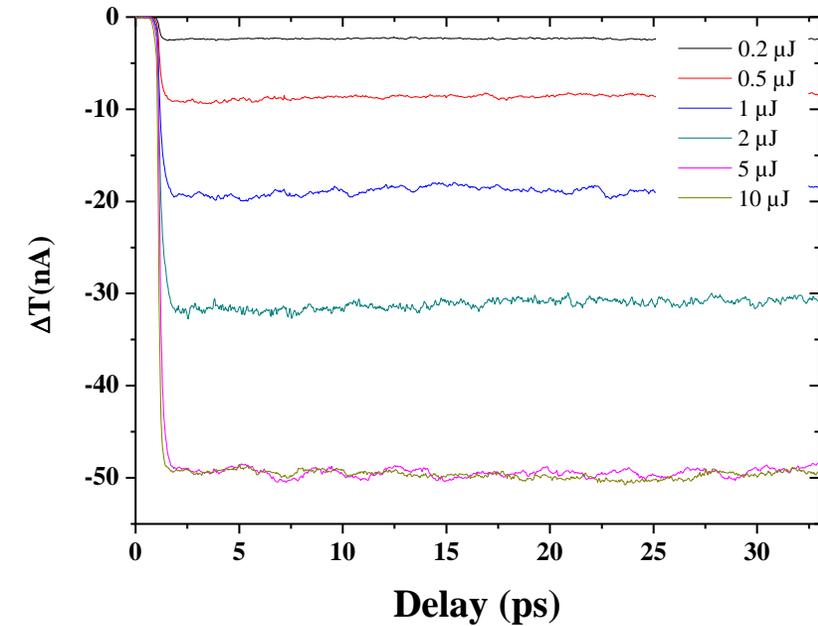
Optical-pump THz probe experiments in transmission at 0°

Sketch of the experimental set-up



Result

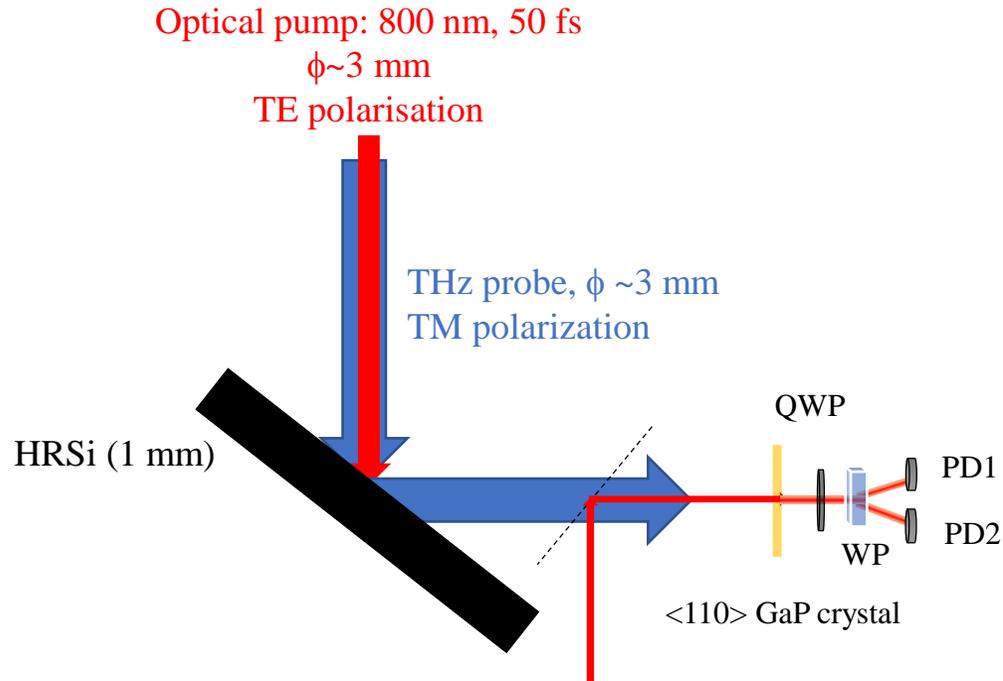
- The THz pulse is sampled at its maximum
- The pump is delayed with respect to the THz pulse



After 1 ps and above 10 μ J of pump the transmission of the THz pulse $< 10\%$!

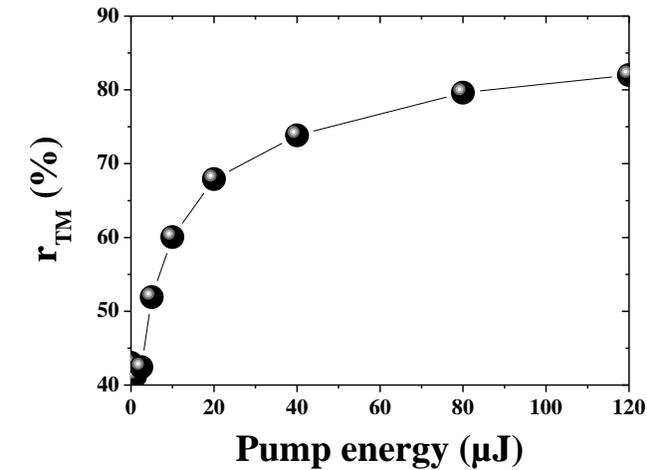
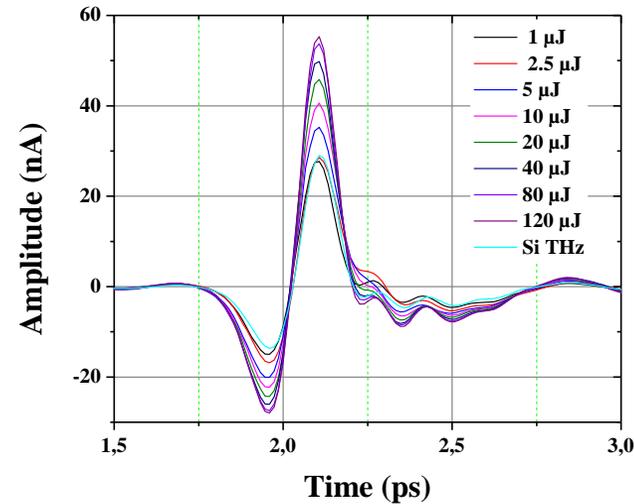
Optical-pump THz probe experiments in reflection at 45°

Sketch of the experimental set-up



Result

- The pump pulse hits HRSi 10 ps before been sensed by the THz pulse



The reflection coefficient for the electric field increases up to 85% !

Model for index variation induced by photo-carriers in HRSi

- The pump pulse is absorbed and generates $N(E)$ carriers in HRSi within the absorption length l_{abs} ($\sim 10 \mu\text{m}$ at 800 nm)

$$N(E) \approx T\alpha \frac{E\lambda}{hc} \frac{4}{\pi l_{abs}\phi^2}$$

- T: transmission coefficient for the pump pulse
- α : absorption of pump within absorption length l_{abs}
- E: energy of the pump pulse
- Φ : beam diameter transmission
- λ : pump wavelength
- c: speed of light
- h: Plank's constant

Example: for $\phi=3 \text{ mm}$ and $\lambda= 800 \text{ nm}$ * $N(1 \mu\text{J}) \sim 2.5 \cdot 10^{22} \text{ m}^{-3}$ * $N(100 \mu\text{J}) \sim 2.5 \cdot 10^{24} \text{ m}^{-3}$

- The generated carriers modify the dielectric constant

$$\varepsilon(\omega) = \varepsilon_{HRSi}(\omega) - \frac{\omega_p^2}{\omega(\omega+i\Gamma)} \quad \text{where} \quad \omega_p^2 = \frac{N(E)e^2}{\varepsilon_0 m^*}$$

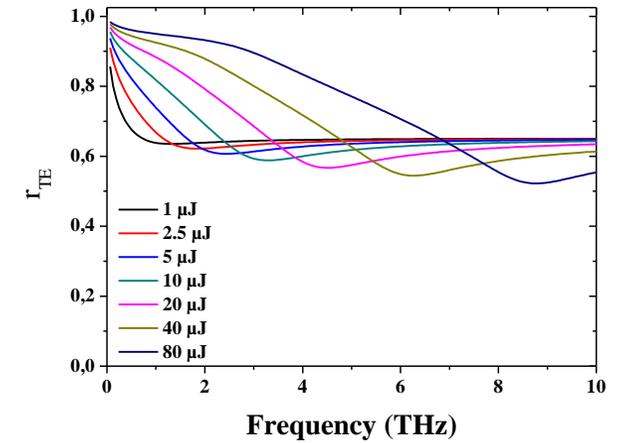
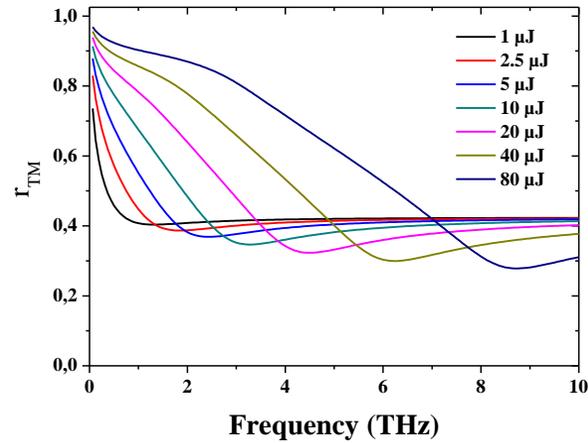
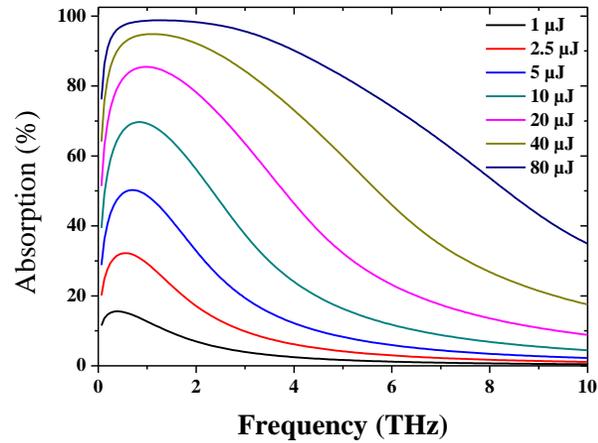
- e and m^* are the charge and reduced mass of the electron, respectively

- In turn, this results in a change of the reflection coefficient and absorption constant

$$r_{TE}(\omega) = \frac{\cos(\theta) - \sqrt{\varepsilon(\omega) - \sin^2(\theta)}}{\cos(\theta) + \sqrt{\varepsilon(\omega) - \sin^2(\theta)}} \quad \text{and} \quad r_{TM}(\omega) = \frac{-\varepsilon(\omega) \cos(\theta) + \sqrt{\varepsilon(\omega) - \sin^2(\theta)}}{\varepsilon(\omega) \cos(\theta) + \sqrt{\varepsilon(\omega) - \sin^2(\theta)}}$$

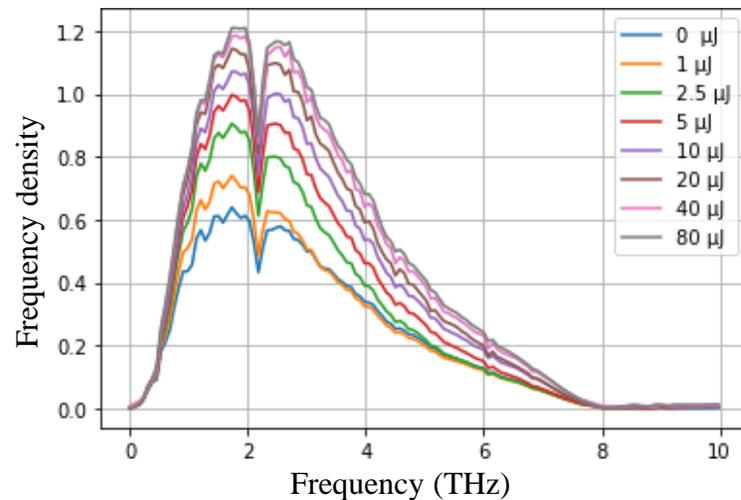
Comparison experiment-theory

- Results of the numerical simulations at $\theta=45^\circ$ of incidence with beam diameter of 3 mm

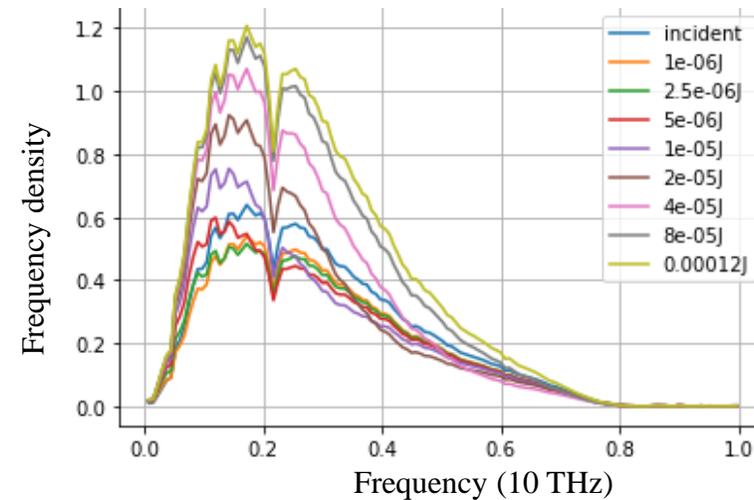


- Comparison with experiment

Experiment (TM polarization)

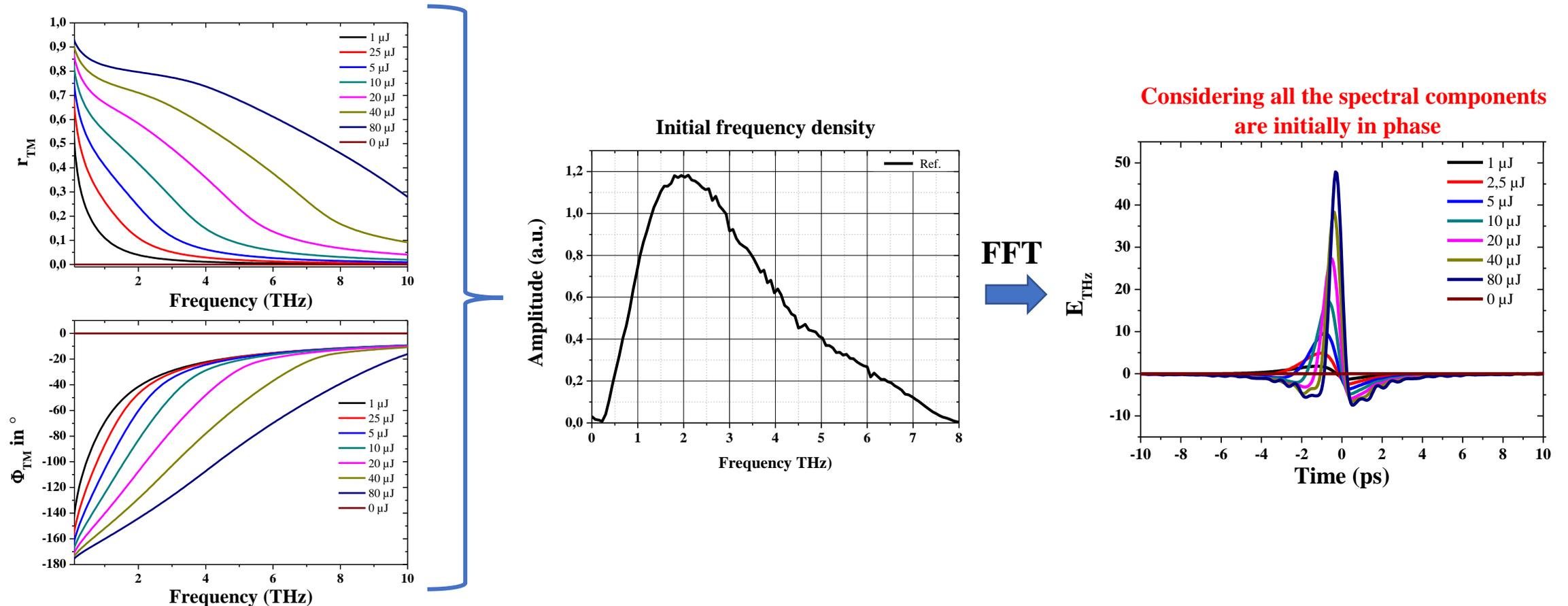


Theory (TM polarization)



Evolution of the reflection coefficient at Brewster's angle

- The pump and THz are collinear, TM polarized, focused on a beam spot of $\sim 7 \text{ mm}^2$ with $\theta_{\text{inc}} \sim 73.4^\circ$ on HRSi



- With 80 μJ pump pulse and within 1 ps, one can modulate r_{TM} from 0 to $\sim 80\%$ in between 0 and 4 THz !
- But since r_{TM} is a complex number, the temporal shape of the THz pulse is expected to be modified

Conclusions

- **We have study the evolution of the transmission and reflection coefficient of a THz pulse on HRSi substrate upon excitation by an amplified femtosecond Ti:Sapphire laser pulse.**
- **We have experimentally shown that the transmission and reflection coefficients can be largely, efficiently and rapidly modulated by the pump pulse.**
- **A simple model where carriers generated by the pump pulse well accounts for the observed phenomena.**
- **This model indicates that at Brewster's angle and for TM polarized optical pump, one will be able to modulate the reflection coefficient of a THz pulse from 0 up to 80% in the 0-4 THz spectral range.**
- **We have shown the temporal shape of the reflected pulse should also be modified.**
- **This study also indicates that HRSi can also be used as ultrafast, broadband and tunable density filter.**